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Thermal Performance Analysis of an Underground Closed Chamber with Internal Heat Sources under Natural Convection

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Abstract: In this article, a combined experimental and numerical study has been performed to 12 investigate the thermal performance of mine refuge chamber (MRC) under natural convection. By 13 14 using heat lamps to simulate human heat loss, a 20-hour heating experiment is carried out in a 15 fifty-person MRC laboratory. A proposed numerical model is validated against the experiment. Furthermore, sensitivity analysis is performed by Fluent software to investigate the effects of 16 17 thermal parameters of rock. Results indicated that: (1) the experimental data and the corresponding numerical prediction have the same trend in air temperature rising with time, and the deviation 18 between the two is less than 10%, which proves that the numerical model is effective; (2) the 19 temperature rise process in a MRC can be divided into air temperature rapid increase stage and air 20 21 temperature slow increase stage; (3) a new analytical method with simplified for predicting air 22 temperature is proposed, it shows that the air temperature growth trend becomes slow with the increase of thermal conductivity, density and specific heat capacity of the rock; (4) the surface heat 23 24 transfer coefficient of vertical walls is the largest and it increases linearly with air temperature.

25 Keywords: Mine refuge chamber; Coal mine safety; Temperature; Natural convection; Surrounding

26 rock; Thermal performance.

Nomenclature						
а	Thermal diffusivity of rock, $a = \lambda/(\rho C_p)$ m ² /h	T'	Fluctuating temperature			
Α	Area, m ²	T_0	Initial rock temperature, \mathcal{C}			
b	Assuming temperature variables, °C	u	Air velocity, m/s			
C ₁	Turbulence model parameter	u'_i, u'_j	Corresponding fluctuating velocity components in the i and j directions			
C ₂	Turbulence model parameter	<i>x</i> , <i>y</i>	Coordinate direction vector			
C ₁ s	Turbulence model parameter	Subscripts				
Ca	Turbulence model parameter	a	Air			
C_{a}	Thermal capacity of air. $J/(kg \cdot K)$	c	Cross section of the original tunnel			
C _n	Thermal capacity of rock, J/(kg·K)	f	flow air			
F_0	Fourier Number	n	Numerical simulation			
0			Elemental directions (i, $i = 1, 2$ and 3			
g	Gravitational acceleration, m ² /s	i, j	according to the rest and redirections)			
			corresponding to the x, y, and z directions)			
G _b	Generation of turbulence kinetic energy due to	u	Unaffected by the external Environment			
	buoyancy (J/s·m ³)	u				
G_k	Generation of turbulence kinetic energy due to the mean velocity gradients $(J/s \cdot m^3)$	W	Walls in a MRC			
G_r	Grashof number, $G_r = g\alpha_v \Delta t l^3 / v^2$	Greek sy	eek symbols			
_	Natural convective heat transfer coefficient		Air volume expansion coefficient			
h	between air and wall, $W/(m^2 \cdot K)$	α_v				
h_0	Initial Natural convective heat transfer coefficient between air and wall, W/(m ² ·K)	β	Coefficient of thermal expansion, 1/K			
k	Turbulent kinetic energy (J/kg)	Δ	Difference			
k_1	Gradient for surface heat transfer coefficient	ε	Turbulent energy dissipation (J/kg·s)			
k_2	Assuming constants for K	ρ	Density, kg/m ³			
Κ	Gradient for air temperature increasing	λ	Thermal conductivity, W/(m·K)			
l	Turbulence length scale, m	μ	Dynamic viscosity, kg/m·s			
<i>m</i> , <i>n</i>	Assuming constants for K	τ	Time, h			
Ν	Number of people in a MRC	$\mu_{ au}$	Turbulent eddy viscosity, (kg/m·s)			
Р	Pressure, Pa	υ	Kinematic viscosity (m ² /s)			
P_r	Planck number, $P_r = u/a$	σ_k	Turbulence model parameter			
q	Heat flux on wall surface, W/m ²	σ_{ε}	Turbulence model parameter			
q_0	Heat generation rate per person, W	Acronym	S			
Q	Total heat generation rate, W	Actonyms				
r	Radius of the surrounding rock, m	MRC	Mine refuge chamber			
r_0	Equivalent radius of the cylinder, m	PCM	Phase change material			
S	Modulus of the mean rate-of-strain tensor	UB	Underground building			
Т	Temperature, °C	Re	Reynolds number			

28 **1. Introduction**

With the rapid development and application of underground energy and space, the safety of underground space is becoming an important issue [1]. According to statistics of underground fire and explosion accidents, only a few people among the deaths died directly from the fire and blast wave, and up to 80% of the people indirectly died from carbon monoxide and suffocation [2-5]. The application of underground emergency refuge system, such as the tunnel rescue station in subway, mine refuge chamber (MRC) and movable shelter in mine, plays an important role in reducing deaths for underground accidents [6-8].

MRC is the most important emergency refuge system for underground mine in China, it is also 36 37 applied in other developed mining countries. A MRC normally consists of a living room and two transition rooms [9], and it needs to serve at least 96 h [10, 11]. Due to the heat generated by human 38 39 metabolism, the air temperature in the MRC may beyond the allowable range of human's thermal tolerance. The recommended apparent temperature in MRC is below 35 °C [12]. Li et al. [13] 40 concluded that human responses change significantly when exposed in an environment with the 41 42 temperature of 33 °C or the relative humidity of 85% in MRC. Du et al. [14] recommend that the temperature and relative humidity in MRC should be less than 31 °C and 80% RH. It needs to be 43 mentioned here that the conventional refrigeration technology can't be applied in MRC, because the 44 45 power may be interrupted after an accident. Therefore, it is imperative to seek new cooling methods for MRC. Jia et al. [15] proposed a temperature control strategy by using ice storage capsule for 46 47 movable MRC. The accuracy of the strategy was verified by a 24-h experiment in a closed cabin. Wang et al. [16] developed an ice thermal storage system, the system was determined in a 48 fifty-person MRC for approximately 64.57 h. Xu et al. [17] proposed a non-electric cooling scheme 49 for placing the encapsulated ice plates directly in the MRC, one plate has an average cooling load of 50 51 14.3 W. Yang et al. [7] designed an open cycle carbon dioxide refrigerator system. A test showed that the system had 1200 W cooling power. Yuan et al. [18] proposed a coupled cooling method and 52 application of phase change material (PCM) combined with pre-cooling of the envelope for MRC, 53 54 the method considered the applicable temperature range of PCM and the cold storage function of the rock. Gao et al. [19, 20, 21] studied the temperature controlling characteristics of PCM plates 55 and PCM seats used in a fifty-person MRC, the coupled heat transfer characteristics of surrounding 56 57 rock, air and PCM were considered in their model.

58 The heat transfer between the heat source, air and surrounding rock in MRC is a dynamic coupling process, thus the calculation of heat transfer is very complex. In recent years, some studies 59 on coupling heat transfer characteristics of air and surrounding rock in the underground building 60 (UB) have been reported. Yuan et al. [22, 23] established a mathematical heat transfer model for 61 62 underground engineering envelope, the model provided a rapid and accurate solution for calculation 63 of heat transfer. Their results indicated the thermal conductivity of the rock is an important factor of 64 the heat transfer. Xiao et al. [24] proposed a method to calculate the transient heat flow through the 65 envelope of an underground cavern and proved that the method has a good agreement with the numerical results. Liu et al. [25] presented a numerical model for the simultaneous heat transfer 66 between air and the tunnel surface. The model was validated against experimental data applied to an 67 underground tunnel. Kajtar et al. [26] presented a dimensioning method for shallow buried UB in 68 69 terms of heat transfer characteristics and thermal comfort. The method was in favor of the quick 70 sizing of the required heating and cooling performance of UB. Szabó et al. [27] developed a new 71 dynamic dimensioning method for shallow buried UB. According to the method, there is no significant change in air and wall temperature after 1000 h, as well as the heat flux through the wall. 72

73 Sasmito et al. [28] studied the thermal management strategies of a dead end ventilated through a 74 pipe in an underground mine, their results showed that several control parameters, such as virgin 75 rock temperature, ventilation temperature and ventilation amount, have a significant effect on air temperature control. Habibi et al. [29] built a ventilation model calibrated against pressure, quantity 76 and temperature results to simulate the airflow and heat conditions for a coal mine. For both flow 77 78 and temperature, the predicted results agreed to within 90% accuracy of the actual measurements. 79 Li et al. [30] pointed out that relative roughness plays an important role in the heat transfer of 80 underground tunnels. As the relative roughness increase, the temperature drop and the cooling efficiency increase gradually. 81

In summary, for the temperature control in MRC, the development of non-powered refrigeration technology has attracted much attention, few studies have focused on the thermal performance of the MRC. In this article, the characteristics of the dynamic coupling heat transfer process between the surrounding rock and the air in a heated MRC are mainly studied. A fifty-person MRC is selected as a study case. A heating experiment is carried out in a MRC laboratory to present the air temperature rising trend. Then ten cases with different parameters are designed to investigate the thermal performance of MRC under natural convection by using Fluent 18.0 software.

89 2. Experimental setup

90 2.1. Experimental environment

91 A MRC is usually located in a deep underground coal mine. Due to the safety needs of the coal 92 mine and the lack of mineral intrinsic safety heating equipment in the market, the heating experiment is conducted in a shallow MRC laboratory. The laboratory can accommodate 50 people 93 in the living room with 20 m in length, 4 m in width and 3 m in height. The top of the living room is 94 0.6 m above the ground. The wall was made of concrete with the density of 1600 kg/m³, the specific 95 heat capacity of 840 J/(kg K) and the thermal conductivity of 0.81 W/(m K). The thickness of the 96 97 vertical and bottoms wall is 0.6 m, 0.4 m for the top wall. A polyurethane insulation layer is covered 98 on the top wall with a thickness of 0.08 m. The thermal conductivity of polyurethane is 0.024 99 W/(m K). In addition, the MRC laboratory is located in a factory, it can avoid the sun shining on the 100 walls of the MRC laboratory.

101 The experiment is performed in September, the heating process starts at 8 a.m. The atmospheric 102 temperature ranges at $22 \sim 26$ °C during the day (from 8 a.m. to 7 p.m.) and $18 \sim 22$ °C in the night.

103 2.2. Measurement and data acquisition

When a man sitting quietly in the MRC, the heating rate can be assumed to be 120 W and the rate of CO_2 released is 0.30~0.35 ml/min [9, 31, 32]. Some measures need to be taken to remove the CO₂ gas. When calcium hydroxide is used to remove CO_2 , the heat load is 20-25 W per person. But the heat load may not be released into MRC through a reasonable design of purification equipment. If CO_2 is removed by fresh air, there will be no heat generated. Therefore, the heat generated by facilities in MRC is not considered in our study.

In the experiment, 40 heat lamps with 150 W, representing the heat production of 50 persons, are divided into 4 rows×10 columns. The row spacing is 1 m and the column spacing is 1.2 m. All heat lamps are 1 m above the bottom as illustrated in Fig. 1. Six measuring points are respectively set at the three horizontal levels of 0.5 m, 1 m, and 1.5 m. The distance from these measuring points to the near side wall is 1 m. The location of the measuring points can be seen in Fig. 1.



- 116
- 117 118

Fig. 1. Distribution of heat lamps and temperature sensors.

The calibrated PT100 (model: WZP-PT100 A; manufacturer: Hangzhou Meacon automation 119 120 technology Co., Itd, China) with a measuring range of $-50 \sim 250$ °C and accuracy of 0.15 °C is chosen to measure the air temperature. And a calibrated infrared thermal imager (model: CEM 121 DT-9868; manufacturer: Shenzhen CEM Co., Itd, China) with a display accuracy of 0.1 °C is used 122 to test the initial temperature of the surrounding rock wall surface. The air temperature is collected 123 124 by a data acquisition subsystem, then transmitted to a temperature monitoring platform and 125 automatically recorded once per minute. The power stability of the heating lamps is guaranteed by a stable voltage power control cabinet. The working condition of the heating lamps and the air 126 temperature measurement system can be controlled in an independent control room. Fig. 2 shows 127 128 the schematic of the experimental apparatus.





Fig. 2. Schematic of the experimental apparatus.

- 132
- 133 2.3. Experimental procedure
- 134 The key steps of the experiment are as follows:
- (1) Check to make sure that make sure that all heating lamps and all temperature sensors canwork properly and the data can be automatically recorded.
- (2) Prior to heating, measure five points on each wall, the average temperature of all measuring
 points is taken as the initial temperature of the wall. The value is 22.3 °C. After the measurement,
 the tester leaves the lab and closes the laboratory's door.
- (3) Prior to heating, turn on the temperature monitoring platform half an hour before heating totest the initial air temperature, the average air temperature is taken as the initial air temperature in

- 142 the MRC laboratory. The value is 25 $^{\circ}$ C.
- 143 (4) Turn on the heating lamps in the control room, heating lasts more than 20 h.
- 144 (5) End the experiment and save the experimental data.
- 145 Fig. 3 shows the actual heating scene in the experiment.



147 148

149

Fig. 3. Normal operation of the heating experiment.

150 **3. Computational details**

151 3.1. Analytical model

The buried depth of an underground mine is usually greater than 100 m. According to [33], the heat transfer characteristics of underground buildings with a buried depth greater than 12 m are not affected by the ground environment temperature. The thermal performance of deep buried underground buildings can be analyzed based on the semi-infinite object heat transfer theory. The controlling equations can be established in the one-dimensional coordinate system.

157 The equation for calculating the changes in the temperature of MRC is:

$$\frac{\partial T(x,\tau)}{\partial \tau} = a \frac{\partial^2 T(x,\tau)}{\partial x^2} \tag{1}$$

$$\begin{cases} \lambda \frac{\partial T(x,\tau)}{\partial \tau} \Big|_{x=0} = h \left(T_f(\tau) - T(x,\tau) \right) \\ \lim_{x \to \infty} \frac{\partial T(x,\tau)}{\partial \tau} = 0 \end{cases}$$
(2)

160

162

158

161 Initial condition can be described as

$$T(x,0) = T_0 \tag{3}$$

In order to simplify the heat transfer process in a heated MRC under the natural convection,several assumptions could be made as follows:

165 (1) The shape of the MRC is cylindrical. Huang et al. [33] proved that the temperature contour

- 166 in the rock formation is approximately circular for different shape tunnels. The equivalent
- 167 radius of the cylinder can be calculated as $r_0 = \sqrt{A_c/\pi}$;
- (2) The heat production rate of each person is equal and constant because people in MRC are
 basically quiet;
- (3) The heat absorbed by air can be ignored because the specific heat capacity of air is much smaller than that of rock;
- (4) Heat transfer on the wall surfaces is uniform and can be regarded as a constant because the
 thermal parameters of the surrounding rock are uniform;
- (5) The temperature inner the surrounding rock is equal everywhere at the initial time because
 the heat transfer characteristics of MRC is not affected by the environment temperature.
- Therefore, in the cylindrical coordinate system, the governing equation of heat conduction can bedescribed as follows

178
$$\frac{\partial T(r,\tau)}{\partial \tau} = a \left(\frac{\partial^2 T(r,\tau)}{\partial r^2} + \frac{1}{r} \frac{\partial T(r,\tau)}{\partial r} \right)$$
(4)

179 The boundary conditions can be described as

180
$$\begin{cases} \left. \lambda \frac{\partial T(r,\tau)}{\partial r} \right|_{r=r_0} = q = \frac{Q}{A_w} = \frac{N \cdot q_0}{A_w} = const\\ \lim_{r \to \infty} T(r,\tau) = T_0 \end{cases}$$
(5)

181 The initial condition is

182

184

188

190

194

$$T(r,0) = T_0 \tag{6}$$

183 An approximate solution regarding the temperature of the wall surface was recommended as [33]

$$T(r_0, \tau) = t_0 + \frac{q \cdot r_0}{\lambda} \frac{1.13\sqrt{F_0}}{1+0.38\sqrt{F_0}} \qquad (F_0 = a \cdot \tau/{r_0}^2)$$
(7)

185 According to Newton's law of cooling

- 186 $q = h[T_f(\tau) T(r_0, \tau)]$ (8)
- 187 Thus, the air temperature in a MRC can be calculated as
 - $T_f(\tau) = T_0 + q \left(\frac{1}{h} + \frac{r_0}{\lambda} \frac{1.13\sqrt{F_0}}{1+0.38\sqrt{F_0}}\right)$ (9)
- 189 The natural convective heat transfer coefficient between the air and wall can be calculated as [21]
 - $h = \frac{\lambda_a \cdot Nu}{l} = \frac{\lambda_a}{l} \times 0.59 \times (G_r \cdot P_r)^{\frac{1}{4}}$ (10)
- 191 The air temperature increase value in the MRC was then computed by
- 192 $\Delta T_f(\tau) = T_f(\tau) T_0 = q \left(\frac{1}{h} + \frac{r_0}{\lambda} \frac{1.13\sqrt{F_0}}{1+0.38\sqrt{F_0}}\right)$ (11)
- 193 The radius of the heat transfer zone of the rock can be calculated as [34]

$$r = 4\sqrt{a \cdot \tau} \tag{12}$$

According to the Eq. (9), it can be seen that the final air temperature in a MRC has a direct 195 relationship with the rock initial temperature. However, according to the Eq. (11), the air 196 temperature increasing value in MRC has nothing to do with the initial rock temperature, it's just 197 198 affected by the thermal parameters of the rock and the heat load on the walls. It is obvious that the 199 increased value is linearly proportional to the heat load on the walls. The heat load is mainly determined by the power of heat sources and the total surface area of the walls, which has nothing 200 to do with the shape and size of the heat sources (miners in the MRC). In addition, as the buried 201 depth of a MRC is much larger than 12 m, the depth will no longer affect the heat transfer 202

characteristics of the MRC. Therefore, in the following numerical study, the burial depth of MRC,
 the initial rock temperature, the human heat loss, as well as the shape and size of human bodies will
 not be considered as the main influencing factors. And the effects of the thermal conductivity and
 the specific heat capacity, as well as the density of the rock on heat transfer characteristics will be
 emphasized.

208 3.2. Computational model

A computational model of MRC is built with the same inner sizes of the MRC laboratory (length 209 ×width×height: 20×3×4 m). The thickness of the wall is 1.5 m. The surface area of a human body 210 model is 2 m^2 . 50 bodies are divided into 4 rows, as shown in Fig. 4. For the two rows adjacent to 211 the two sides of the room, each row has 13 bodies, the back of the body is 0.3 m from the wall. For 212 the two rows in the middle, each row has 12 bodies, The distance between two bodies' backs is 0.4 213 214 m. The center distance between the two adjacent bodies in a row is 1 m. Meanwhile, in order to obtain a high-quality boundary layer grid, the bottom surface of human body is above the bottom 215 216 0.35 m. 217



218 219

220

Fig. 4. Geometric model of the fifty-person MRC.

The computational grids are generated by software ANSYS ICEM 18.0. Six grids (with a number of grids as 10.2×10^5 , 13.8×10^5 , 17.6×10^5 , 27.5×10^5 , 35.0×10^5 , 41.4×10^5 , respectively) are tested to ensure that the solver and numerical schemes implementation yield results are independent from the grid, as shown in Fig. 5.



Fig. 5. Comparison of air temperature at three different time under six different grids.

It has been concluded that the mesh with 17.6×10^5 grids is sufficient. The maximum grid size of the inner wall surface is 0.1 m. In the fluid zone. 4 prism layers are created along the surrounding. The maximum grid size of human body surfaces is 0.06 m. The maximum grid size of the fluid zone and the solid zone is 0.3 m and 0.5 m, respectively.

232

233 3.3. Numerical methodology

Human bodies are defined as solid zones with a constant temperature of $37 \,^{\circ}$ C. Shadow surfaces will be generated automatically at the surfaces of human bodies in the Fluent software. The surfaces of the human bodies adjacent to the fluid zone are defined as constant heat flux boundary with 60 W/m², and the shadow surfaces adjacent to the human body solid zones are defined as heat flux boundary with 0 W/m². The inner walls of the surrounding rock are defined as the coupled boundary. The outer walls of the surrounding rock are defined as heat flux boundary with 0 W/m².

In combination with the thermal parameters of the surrounding rock of the MRC laboratory and the thermal parameters of common rocks in mine, ten different cases are designed, see Table. 1.

Table. 1 Thermal	physical	parameters for	the ten	numerical	cases
------------------	----------	----------------	---------	-----------	-------

NO	$T_{f}(0)$	T ₀	λ	ρ	C _p	$ au_{ m u}$	$ au_{\mathrm{n}}$
NO	°C	°C	W/(m K)	kg/m ³	J/(kg K)	h	h
1	25	22.3	0.81	1600	840	64.81	60
2	20	20	1	2400	920	86.25	60
3	20	20	1.50	2400	920	57.50	20
4	20	20	2	2400	920	43.12	20
5	20	20	2.50	2400	920	34.50	20
6	20	20	3	2400	920	28.75	20
7	20	20	2	2400	800	37.5	20
8	20	20	2	2400	1100	51.56	20
9	20	20	2	2000	920	35.93	20
10	20	20	2	1500	920	26.95	20

The parameters in the NO.1 case correspond to the experiment. The remaining cases are designed to investigate the effects of thermal parameters of the rock, the initial temperature of the air and the rock are set as 20 °C. According to Eq. (12), the time that the heat transfer characteristics of the MRC model unaffected by the external environment for each case is different, the unaffected time (τ_u) for each case is shown in Table. 2. The numerical simulation time (τ_n) for NO.1 and NO.2 is 60 h, for the remaining cases, τ_n is 20 h.

251 3.4 Turbulence model

The airflow velocity near the walls in the MRC induced by buoyancy is estimated as $0.02 \sim 0.3$ m/s, the Reynolds number (Re) is calculated as $0.22 \times 10^5 \sim 1.08 \times 10^5$. Therefore, the air flow in the MRC is considered to be turbulent.

255 The effects of turbulence are modeled frequently by using the three models of Standard $k-\varepsilon$, RNG k- ε and Realizable k- ε [35]. Wu et al. [36] proved that, for conjugate turbulent natural 256 257 convection in a differentially heated cavity, the three models were acceptable in terms of the performance of predicting the time-averaged quantities, and the variation between them was very 258 small. Franke et al. [37] indicated that realizable $k-\varepsilon$ turbulence model had a general good 259 performance for wind flow around buildings. Sørensen et al. [38] indicated that the Realizable k-e 260 261 had an overall good performance for indoor air flow. Through a study of the natural convection 262 phenomena inside a wall solar chimney, Bacharoudis et al. [39] proved that the realizable $k-\varepsilon$ model was likely to provide superior performance for flows boundary layers under strong adverse 263 pressure gradients. Piña-Ortiz et al. [40, 41] proved that the Realizable k- ε model was the best 264 model for natural convection in a cubic cavity with the lowest temperature difference. Therefore, 265 the Realizable $k-\varepsilon$ model is selected for the current study. 266

The buoyancy-induced turbulent air flow within the MRC is governed by the following unsteadyReynolds Averaged Navier–Stokes equations [36]:

269 The continuity equation is

270

275

$$\frac{\partial \rho_a}{\partial \tau} + \frac{\partial (\rho_a \cdot u_i)}{\partial x_i} = 0 \tag{13}$$

271 The momentum equation is

272
$$\frac{\partial u_i}{\partial \tau} + \frac{\partial (u_i \cdot u_j)}{\partial x_j} = -\frac{1}{\rho_a} \frac{\partial P}{\partial x_i} + \frac{1}{\rho_a} \frac{\partial}{\partial x_j} \left[\mu \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) - \rho_a \overline{u'_i \cdot u'_j} \right] - g_i \beta (T - T_0)$$
(14)

The Boussinesq approach is applied for the effect of gravity force, the energy equation with Boussinesq assumption is

$$\frac{\partial T}{\partial \tau} + \frac{\partial (u_j T)}{\partial x_j} = \frac{1}{\rho_a} \frac{\partial}{\partial x_j} \left[\mu \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) - \rho_a \overline{u'_i \cdot T'} \right]$$
(15)

276 The realizable $k-\varepsilon$ model consists of the following two transport equations[42]:

277
$$\frac{\partial}{\partial \tau} (\rho_a k) + \frac{\partial (\rho_a k u_j)}{\partial x_j} = \frac{\partial}{\partial x_j} \left[\left(\mu + \frac{\mu_\tau}{\sigma_k} \right) \frac{\partial k}{\partial x_j} \right] + G_k + G_b - \rho_a \varepsilon$$
(16)

278
$$\frac{\partial}{\partial \tau}(\rho_a \varepsilon) + \frac{\partial(\rho_a \varepsilon u_j)}{\partial x_j} = \frac{\partial}{\partial x_j} \left[\left(\mu + \frac{\mu_\tau}{\sigma_\varepsilon} \right) \frac{\partial \varepsilon}{\partial x_j} \right] + \rho_a C_1 S \varepsilon + \rho_a C_2 \frac{\varepsilon^2}{k + \sqrt{\nu \cdot \varepsilon}} - \rho_a C_{1\varepsilon} \frac{\varepsilon}{k} C_{3\varepsilon} G_b \tag{17}$$

279 3.5 other settings

The gravity value is 9.81 m/s^2 . The operating pressure is 102325 Pa. The air operating density is 1.225 kg/m³. The enhanced wall treatment with pressure gradient effects and thermal effects are taken into account, as well as the full buoyancy effect. Pressure-implicit with splitting of operators (PISO) is used for the pressure-velocity coupling. The pressure is discretized by using the body force weighted schemes. The energy and momentum are discretized by using the second-order upwind schemes. The convergence absolute criteria for energy is set to 10^{-6} , for other items is 10^{-3} . The time step is 10 s.

287 4. Results and discussion

288 4.1. Model validation

Fig. 6 plots the variation of average air temperature with heating time at three different levels, i.e.0.5 m, 1 m, 1.5 m.

It is observed from Fig. 6 that the average air temperature at the three levels has the same growth trend, the air temperature monotonically increases over time. However, there are slight differences in temperature at different height levels, indicating that the temperature increases with the height. During the period from 2 to 10 h, the air temperature at 1 m level is $0.3 \sim 0.4$ °C higher than that at 0.5 m level, but $0.2 \sim 0.4$ °C lower than that at 1.5 m level.

296 297



298 299

Fig. 6. Variation of air temperature at three different levels, i.e. 0.5 m, 1 m, 1.5 m, with heating time

300

Fig. 7 shows the comparison of average air temperature between the experimental data and the numerical results of NO.1 case.



306

Fig. 7. Comparison of numerical results and experimental data.

It can be seen from Fig. 7 that the average air temperature obtained through the experiment and numerical method has the same trend of rising monotonically with time. In the experiment, the average air temperature in the MRC rises from 25 % to about 29.5 % in less than 0.5 h after the heating lamps work. Afterward, the rising trend gradually becomes slow over time. It can be found that the predicted air temperature is higher than the experimental one.

312 According to Eq. (12), the $\tau_{\rm u}$ value is 10.37 h. From 0.5 h to 10.3 h, the average air temperature difference between the experimental value and the predicted value is 0.6~0.8 °C. The temperature 313 deviation between the experimental temperature and the predicted temperature ranges from 8.5% to 314 14.5%, referencing to the initial air temperature (25 °C) in the MRC. When taking the initial 315 surrounding rock temperature (22.3 $^{\circ}$ C) as the reference, the deviation value is 6.3%~9.5%. The 316 317 temperature difference is mainly attributed to two aspects: firstly, the location of the measuring 318 points is relatively low, the measuring value may be smaller than the actual value; secondly, the MRC laboratory experienced a period of time prior to heating, the internal temperature of the rock 319 may not be uniform. From 10.3 h to 20 h, the air temperature rising trend in the experiment 320 becomes slower, and the temperature difference and the deviation both increase with the heating 321 322 time. It could be explained by the fact that after about 10 h of heating, the experiment began to be affected by the external environment. 323

At the beginning of the heating, since the experimental result shows that the air temperature is very sensitive and the initial air temperature is higher than the rock temperature, it is not appropriate to take the initial air temperature as the reference. According to Eq. (9), the rock initial temperature is a reasonable reference. So during the unaffected time of the MRC laboratory, the deviation value is less than 10%. On the other hand, it can be easily found that both the numerical results and the experimental data have an obviously similar trend in air temperature increases. Therefore, it can be concluded that the numerical model is effective.

331 4.2. Air temperature distribution in the MRC

Fig.8 shows the temperature contours of the center cross-sectional at different time.



Fig. 8. Contours of temperature distribution at different time.

336 It can be seen from Fig. 8 that the air temperature in MRC is not uniform. The air temperature 337 above the top surface of the human body is higher than that below the top. In the above part, the air temperature decreases with the height. In the below part, the air temperature increase with the height. 338 The air temperature difference between the top and bottom increases with the heating time. 339

340 4.3. Trend of air temperature rising in the heating process

Fig. 9 and Fig. 10 demonstrate the variation of average air temperature with heating time (τ) and 341 the square root of heating time ($\sqrt{\tau}$) in 60 h for NO.1 case, respectively. 342





It can be seen from Fig. 9 that the air average temperature in the MRC monotonically increases with time. At the beginning of heating, the air temperature in the room rises quickly, the temperature increasing from 25 $^{\circ}$ C to 30.2 $^{\circ}$ C only experiences 0.35 h. Afterward, the increasing rate gradually slows down, the air temperature rising to 35 $^{\circ}$ C takes about 20 h.





351 352

It can be seen from Fig. 10 that the air temperature rising trend approximately exhibits two linear growth stages. The gradient of air temperature rising in $\sqrt{\tau} < 0.5$ h is obviously larger than that in $\sqrt{\tau} > 0.5$ h. Therefore, the process of the air temperature rising in the heated MRC can be divided into two stages, they are air temperature rapid increase stage and air temperature slow increase stage.

358 During the air temperature rapid increase stage, assume that the air temperature is evenly 359 distributed, according to the principle of energy conservation, there is

360
$$Qd\tau = mC_a dT_f(\tau) + hA_w [T_f(\tau) - T(r_0, \tau)] d\tau$$
(18)

The time of the air temperature rapid increase stage is short (less than 0.5 h), during this time, the temperature of the rock surface changes less. Therefore, it can be assumed that the temperature of the walls does not change. If the initial air temperature is equal to the initial rock temperature, that is

$$T(r_0, 0) = T_f(0) = T_0 \tag{19}$$

366 The Eq. (18) can be solved as

367

365

$$T_f(\tau) = \frac{q}{h} \left(1 - e^{-\frac{hA_W}{m \cdot C_a} \tau} \right) + T_0$$
⁽²⁰⁾

During the air temperature slow increase stage, it can be easily found from Fig. 10 that the air temperature is obviously linearly related to the square root of heating time, that is

370
$$T_f(\tau) = K\sqrt{\tau} + b = q \times f(\lambda, \rho, C_{P_i})\sqrt{\tau} + \frac{q}{h} + T_0$$
(21)

371 4.4. Convective heat transfer coefficient on wall of enclosure structure

Fig. 11 shows the surface heat transfer coefficient of the bottom, the top, as well as the verticalwalls changes with the average air temperature in the MRC.





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377

Fig. 11. Wall convective heat transfer coefficient varies with temperature.

378 It can be seen from Fig. 11 that the surface heat transfer coefficient value in different directions are not equal. The value on the vertical wall is the largest and on the bottom is the smallest. The 379 value of the vertical wall increase monotonically linearly with air temperature, but for the bottom 380 and the top wall, it does not change substantially. It means that, for MRC with same space volume, 381 the trend of air temperature rising can be slowed down by increasing the surface area of the vertical 382 walls. The predicted average natural convection heat transfer coefficient is $3.9 \sim 4.8 \text{ W/m}^2 \text{ K}$ in the 383 31~39 °C environment. Yoon et al. [43] found that the average natural convection heat transfer 384 coefficient is 4.53 W/m² K, by performing a test in an underground tunnel on a summer day with 385 atmospheric temperature range from 23.84 °C to 29.47 °C. But in their test, the sampling points are 386 located on the both-side vertical walls, and the effect of the wall roughness on the value is not well 387 estimated. To some extent, the predicted value is close to the test one. 388

389 For the average surface heat transfer coefficient of the MRC, the value shows a linear increase

390 with the air temperature, that is

391

$$h(\tau) = k_1 [T_f(\tau) - T_f(0)] + h_0$$
(22)

The difference in the convection heat transfer coefficient is mainly due to the uneven distribution of air velocity in the MRC. Fig.12 shows the air velocity magnitude distribution in the MRC at 40 h.

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396 397 398

According to Fig. 12, as far as the air velocity direction is concerned, the air moves upward around the heat source bodies, moves horizontally both near the top surface and the bottom surface, moves downward along the vertical wall surfaces. Regarding the air velocity magnitude, the wind speed near the vertical wall is the largest, near the bottom wall is smallest. The maximum wind speed near the vertical wall is 0.25 m/s, 0.12 m/s near the top, and 0.02 m/s near the bottom.

404 *4.5.* Effect of thermal parameters of surrounding rock

Fig. 13 plots the air temperature increase with time for the nine different cases NO.2~NO.10.

It can be seen from Fig. 13 that the air temperature monotonically increases with time. At the 406 initial stage of the heating, the air temperature rapidly rises to 27.8~28.3 °C from 20 °C in less than 407 0.5 h, the curves of the air temperature basically coincide in the nine different cases, it indicates that, 408 409 during the air temperature rapid increase stage, the λ , ρ and C_P of surrounding rock have no obvious effect on the air temperature rise, which is in good agreement with Eq. (20). When the 410 411 heating time $\tau > 0.5$ h, the air temperature rising trend with time gradually becomes slow, and the 412 greater the λ , ρ and C_P of the rock, the slower the air temperature rises. It can be concluded that, 413 during the air temperature slow increasing stage, the air temperature rise rate decreases as the λ , ρ and C_P of the rock increase. 414

















Fig. 14. Variation of average air temperature with $\sqrt{\tau}$ at different λ , ρ and C_P .

It can be seen from Fig. 14 that, when $\sqrt{\tau} > 1.5$ h, the average air temperature increases linearly with $\sqrt{\tau}$. It is not difficult to judge that the value of λ , ρ and C_P mainly influences the linear

growth gradient (K), the K value decreases as the λ , ρ and C_P of the rock increase.

4.6. Accurate prediction at the air temperature slow increase stage

(1) Establish the new analytical method

Fitting the original data in Fig. 9 and Fig. 13, the linear fitting formula for different λ , ρ and C_P of the rock is obtained, as shown in Table. 2.

Table. 2 The fitting relation for different λ , C_P and ρ of the rock.							
λ	λ $ ho$ $C_{ m p}$		Linger fitting formula	V	\mathbf{D}^2		
W/(m·K)	kg/m ³	J/(kg·K)	- Linear fitting formula	Λ	ĸ		
0.81	1600	840	y = 1.3014x + 29.772	1.3014	0.9999		
1	2400	920	y = 0.9082x + 27.688	0.9082	1		
1.5	2400	920	y = 0.7349x + 27.652	0.7349	1		
2	2400	920	y = 0.6323x + 27.645	0.6323	0.9999		
2.5	2400	920	y = 0.5601x + 27.643	0.5601	0.9999		
3	2400	920	y = 0.5094x + 27.636	0.5094	0.9998		
2	1500	920	y = 0.7915x + 27.639	0.7915	0.9999		
2	2000	920	y = 0.6842x + 27.657	0.6842	0.9996		
2	2400	920	y = 0.6323x + 27.645	0.6323	0.9999		
2	2400	800	y = 0.6764x + 27.645	0.6764	0.9998		
2	2400	1100	y = 0.5768x + 27.651	0.5768	1		

Fig. 15 plots the K value varies with $1/\sqrt{\lambda}$ (value of ρ and C_P are same) and $1/\sqrt{\rho \cdot C_P} \times 10^3$ (λ value is same), respectively, as well as the corresponding fitting line.



437 It can be seen from Fig. 15 that the *K* value has a linear relationship with $1/\sqrt{\lambda}$. Therefore, the 438 Eq. (21) can be further expressed as follow:

$$T_f(\tau) = q\left(\frac{1}{\sqrt{\lambda}} - 0.036\right) \times f(\rho, C_P)\sqrt{\tau} + \frac{q}{h} + T_0$$
(23)

440 It can be also seen from Fig. 15 that the *K* value has a linear relationship with $1/\sqrt{\rho \cdot C_P} \times 10^3$. 441 Therefore, the Eq. (21) can be further expressed as follow:

$$T_f(\tau) = qf(\lambda) \left(\frac{1}{\sqrt{\rho C_P}} \times 10^3 + 0.024\right) \sqrt{\tau} + \frac{q}{h} + T_0$$
(24)

443 According to Eq. (23) and Eq. (24), *K* can be expressed as:

$$K = qf(\lambda, \rho, C_P) = k_2 \left(\frac{1}{\sqrt{\lambda}} + m\right) \left(\frac{1}{\sqrt{\rho C_P}} \times 10^3 + n\right)$$
(25)

Taking the thermal parameters in these ten cases and the corresponding *K* values into Eq. (25), it can be solved that $k_2 = 1.4001$, m = -0.0491, n = 0.0109. So *K* can be expressed as:

447
$$K = q \times f(\lambda, \rho, C_{P_{\lambda}}) = 1.4 \left(\frac{1}{\sqrt{\lambda}} - 0.05\right) \left(\frac{1}{\sqrt{\rho C_{P}}} \times 10^{3} + 0.01\right)$$
(26)

448 Taking Eq. (26) and $q = \frac{6000}{2 \times (3+4) \times 20 + 2 \times 3 \times 4} \approx 19.74$ into Eq. (21), then converting the unit of τ

449 from second (s) to hour (h), there is

450
$$T_f(\tau) = 1.18q \left(\frac{1}{\sqrt{\lambda}} - 0.05\right) \left(\frac{1}{\sqrt{\rho \cdot C_P}} + 1 \times 10^{-5}\right) \sqrt{\tau} + \frac{q}{h} + T_0$$
(27)

451 (2) Applicability analysis of the new analytical method

- 452 Fig. 16 shows the air temperature changes over time though three methods for NO.1 and NO.2.
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Fig. 16. Comparison of the curves obtained by three different methods for NO.1 and NO.2

It can be found in Fig. 16 that, for case NO.1 and NO.2, the air temperature calculated by Eq. (27) is closer to the numerical results than the result of Eq. (9). At 60 h, the temperature difference between the value calculated by Eq. (27) and the numerical result is less than 0.2 \degree for NO.1, and less than 0.1 \degree for NO.2. At 100 h, the temperature value calculated by Eq. (27) is 1.03 \degree higher than the value calculated by Eq. (9) for NO.1. Taking the initial rock temperature as a reference, the difference ratio of temperature is 5.03% for NO.1 and less than 3.5% for NO.2.

Regarding the prediction of air temperature in a heated MRC under natural convection, the existing analytical method has a relatively slow temperature growth trend. The new analytical method presented in this paper are closer to numerical results than the existing methods. The difference ratio between air temperature calculated by the proposed method and the existing analytical method is less than 5% during 96 h. The proposed method is more simple and clear than the existing method in terms of expression.

468 **5.** Conclusions

In this study, a heating experiment is conducted and a corresponding numerical case are performed. Furtherly, another nine numerical cases with different thermal physical parameters of rock are designed to study the effect of heat conductivity, density and specific heat capacity of the rock on the thermal performance of MRC. According to the results, the following conclusions can be drawn:

474 (1)The experimental data and the corresponding numerical case results have similar air475 temperature increasing trend, which effectively validates the numerical model.

476 (2) The process of the air temperature increase in MRC under natural convection is divided into 477 air temperature rapid increase stage and air temperature slow increase stage. During the previous 478 stage, air temperature is nearly unaffected by the λ , ρ and C_P of the rock. During the later stage, 479 air temperature growth trend becomes slow with the increase of λ , ρ and C_P of the rock.

(3) The surface heat transfer coefficient of the vertical wall is the largest, and it shows an obvious
linear growth trend with the temperature. The predicted average natural convection heat transfer coefficient
is close to a test result

(4) A new analytical calculation method for predicting the air temperature in a heated MRC under
natural convection is proposed. During 96 h, the difference ratio of temperature predicted by the
proposed method and the existing method is less than 5%.

The significance of the results is that it doesn't just show that, for MRC with same space volume, the trend of air temperature rising can be slowed down by increasing the surface area of the vertical walls. More importantly, through the proposed analytical method, it can be predicted in advance whether the air temperature in a MRC will exceed the allowed temperature during the 96-hour service time under natural convection, to determine whether it needs to take cooling measures in the MRC. Furthermore, it also indicates that the decoration of MRC walls by using decorative sheets with small thermal conductivity is not conducive to the heat dissipation of surrounding rocks.

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