

Reaction Rate Sensitivity of the γ -Process Path

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The location of the $(\gamma,p)/(\gamma,n)$ and $(\gamma,\alpha)/(\gamma,n)$ line at γ -process temperatures is discussed, using recently published reaction rates based on global Hauser-Feshbach calculations. The results can directly be compared to previously published, classic γ -process discussions. The nuclei exhibiting the largest sensitivity to uncertainties in nuclear structure and reaction parameters are specified.

1. INTRODUCTION

Many proton-rich isotopes of naturally occurring stable nuclei cannot be produced by neutron captures along the line of stability. The currently most favored production mechanism for those p-isotopes is photodisintegration of intermediate and heavy elements at high temperatures in late evolution stages of massive stars, the so-called γ -process [1]. Recent investigations have shown that there still are considerable uncertainties in the description of nuclear properties governing the relevant photodisintegration rates. This has triggered a number of experimental efforts to directly or indirectly determine reaction rates and nuclear properties for the γ -process (see, e.g., [2,3,4,5] and references therein). However, many such investigations focussed on nuclei in the γ -process path without considering whether the rates involving these nuclei actually exhibit large uncertainties. In this work the sensitivity of the location of the γ -process path on reaction rates is investigated, showing which nuclei should be preferred in experimental studies.

2. CALCULATIONS

This work makes use of the standard rate set (based on FRDM) [6] which is also used in many stellar models, e.g. [7]. Similar to Table 2 in [1] for $T_9 = 2.5$, the resulting branching points in the photodisintegration path are shown in the last 3 columns of Table 1, for three temperatures $T_9 = 2.0, 2.5, 3.0$. Following the definition in [1], for each isotopic chain with charge number Z the neutron number N is specified at which the condition $\lambda_{\gamma p} + \lambda_{\gamma \alpha} > \lambda_{\gamma n}$ is fulfilled for the first time when following an isotopic chain towards decreasing N . The branching type is indicated by subscripts.

Usually, experimental investigations primarily focus on nuclei close to these branch points. However, they should rather focus on rates which are sensitive to the nuclear

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input, i.e. nuclei for which $\lambda_{\gamma n}$, $\lambda_{\gamma p}$, and $\lambda_{\gamma\alpha}$ are close. To this end, Table 1 also shows the nuclei for which $\lambda_{\gamma p}$ and $\lambda_{\gamma\alpha}$ are within factors $f \leq 3$ or $3 < f \leq 10$, respectively, of the $\lambda_{\gamma n}$ rate. Subscripts indicate which rate is close to $\lambda_{\gamma n}$. Two subscripts indicate that $\lambda_{\gamma p}$ or $\lambda_{\gamma\alpha}$ are within the quoted range but that they are also within a factor of 3 of each other. Note that according to the definition of the factors a nucleus is not repetitively given in the factor 10 columns when it has already been included in the factor 3 column at the same temperature.

3. DISCUSSION

A direct comparison with Table 2 of [1] shows remarkable agreement with a few exceptions. This is surprising insofar as the previous rate predictions made use of a number of simplifying assumptions, such as equivalent square well potentials in the particle channels and total neglect of excited states. The exceptions are Ba, W, Au, Hg where the new branching points are shifted by 2 units to the more neutron-rich side, Pb which is shifted by one unit, and Ce, Gd, Ho, which have become more neutron-deficient by 2 neutrons. Only the branching in Tl has been shifted by a larger amount, the branching point has 4 neutrons less than previously. The branching type was impacted even less: a combined $\gamma p + \gamma\alpha$ branching was changed into a pure $\gamma\alpha$ one in Ba and Au, and a γp one has become a combined $\gamma p + \gamma\alpha$ branching in Ta. Incidentally, almost all impacted nuclei are within the mass range $125 \leq A \leq 150$ and $168 \leq A \leq 200$ where γ -process nucleosynthesis consistent with solar p-abundances was found using the new rates [7], thus underlining the improvement of the rate predictions.

As pointed out above, experiments targeting the sensitive rates given in Table 1 will have direct impact on γ -process nucleosynthesis. Among them, sensitive rates at branching points (coinciding with the nuclei given in the last 3 columns) will be the most important.

In recent investigations checking theoretical rates against newly measured ones it has become apparent that the largest problem is in determining optical α +nucleus potentials at low energies (see [3,4,5] and references therein). Thus, the $\lambda_{\gamma\alpha}$ rates bear the largest inherent uncertainty whereas $\lambda_{\gamma n}$ and $\lambda_{\gamma p}$ have been found generally well predicted, with a few exceptions [4]. A more detailed study of the branching point sensitivity, also exploring a possible modification of γ -process nucleosynthesis and the impact of different α +nucleus potentials, will be published elsewhere in an extended paper.

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Table 1: Branch points and nuclei with large rate uncertainties (see text). The subscripts at each neutron number indicate which rate ($\lambda_{\gamma p}$ or $\lambda_{\gamma\alpha}$) is close to the $\lambda_{\gamma n}$ rate within the given factor, or the branching type, respectively.

Z	\pm factor 3			\pm factor 10			branch point		
	2.0	2.5	3.0	2.0	2.5	3.0	2.0	2.5	3.0
34				42_α			40_α	40_α	$40_{p,\alpha}$
35	46_p	46_p				46_p	46_p	44_p	44_p
36	$44_{p,\alpha}$	44_p		41_p	41_p	44_p	$44_{p,\alpha}$	42_p	42_p
37		48_p	$45_p, 48_p$		45_p		48_p	48_p	46_p
38	43_p	$43_p, 46_p$	46_p	46_p		43_p	46_p	46_p	44_p
39	49_p	49_p	49_p				50_p	50_p	50_p
40	47_p	50_p	50_p		47_p	47_p	50_p	50_p	48_p
41		46_p					50_p	50_p	50_p
42	52_α				52_α	49_p	52_α	50_p	50_p
43		54_p		54_p		$51_p, 54_p$	54_p	52_p	52_p
44	$51_p, 54_\alpha$	51_p	$52_{p,\alpha}$		52_α	51_p	54_α	52_α	$52_{p,\alpha}$
45				55_p		$53_p, 56_p$	56_p	56_p	56_p
46		53_α		$53_\alpha, 56_\alpha$	56_α	$53_{p,\alpha}, 54_{p,\alpha}$	56_α	54_α	$54_{p,\alpha}$
47	$57_p, 60_p$				$57_p, 60_p$	$58_p, 60_p$	58_p	58_p	58_p
48		$55_{p,\alpha}, 58_\alpha$	55_p			$54_\alpha, 58_\alpha$	58_α	58_α	56_p
49		$59_p, 62_p$	$59_p, 62_p$	59_p			62_p	62_p	60_p
50	$59_{p,\alpha}, 62_\alpha$				59_p	$59_p, 60_p$	62_α	$60_{p,\alpha}$	60_p
51	62_α	68_p	$63_p, 68_p$	65_p	63_p		68_p	68_p	66_p
52	$65_\alpha, 70_\alpha$		68_α		68_α	63_α	68_α	68_α	66_α
53	67_p	67_p		72_p		$67_p, 70_p$	70_p	70_p	70_p
54	67_α	70_α	$68_{p,\alpha}$	72_α		$65_\alpha, 70_{p,\alpha}$	70_α	68_α	$68_{p,\alpha}$
55	71_p	74_p	74_p	76_p	71_p	69_p	74_p	74_p	72_p
56	69_α		$72_{p,\alpha}$		$72_\alpha, 74_\alpha$	67_p	74_α	72_α	$70_{p,\alpha}$
57	$73_p, 78_p$	$73_p, 78_p$	78_p			$73_p, 76_p$	78_p	76_p	76_p
58	78_α	74_α	$74_{p,\alpha}$	76_α	76_α	$71_p, 72_{p,\alpha}$	76_α	74_α	$72_{p,\alpha}$
59	77_p		$75_p, 80_p$	$82_p, 84_\alpha$	$75_p, 77_p, 80_p, 82_p$	$77_p, 82_p$	80_p	80_p	80_p
60	80_α	$78_{p,\alpha}$	$73_p, 76_{p,\alpha}, 78_p$	$75_\alpha, 84_\alpha$	$73_{p,\alpha}, 75_\alpha$		78_α	$78_{p,\alpha}$	74_p
61	81_p		79_p	84_α	$77_p, 79_p, 77_p, 81_p$		84_α	82_p	82_p
62	$79_\alpha, 82_\alpha$	77_α		86_α	$79_p, 81_p$	$77_{p,\alpha}, 79_p, 80_p$	84_α	80_p	80_p
63	88_α			86_α	84_α	84_α	88_α	84_α	82_p
64	$85_\alpha, 88_\alpha$	$79_p, 81_p, 86_\alpha$	$79_p, 81_p, 85_\alpha$	79_p		77_α	88_α	84_α	82_p
65	87_α	88_α	$86_{p,\alpha}$	90_α	86_α	$83_p, 88_{p,\alpha}$	88_α	86_α	$84_{p,\alpha}$

Table 1: (Continued)

Z	\pm factor 3			\pm factor 10			branch point		
	2.0	2.5	3.0	2.0	2.5	3.0	2.0	2.5	3.0
66	83_α		$85_\alpha, 86_\alpha$	$87_\alpha, 90_\alpha$	$87_\alpha, 88_\alpha$	88_α	90_α	88_α	86_α
67	$90_{p,\alpha},$ $92_{p,\alpha}$	$83_{p,\alpha},$ $87_{p,\alpha}$	85_α		90_p	$83_p, 87_p$	88_α	88_p	88_p
68	89_α	83_α	$83_p, 90_\alpha$	$91_\alpha, 94_\alpha$	$87_\alpha, 90_\alpha,$ 92_α	$87_\alpha, 88_\alpha$	92_α	90_α	88_α
69	$91_\alpha, 96_\alpha$	$89_p, 94_{p,\alpha}$	$89_p, 94_p$	89_α			96_α	92_p	92_p
70	91_α	$89_\alpha, 94_\alpha$	$92_{p,\alpha}$	$93_\alpha, 98_\alpha$		$87_\alpha, 89_\alpha,$ 94_α	96_α	94_α	$92_{p,\alpha}$
71	$98_\alpha, 100_\alpha$	$93_p, 96_p$	$93_p, 96_p$	95_α	91_p	91_p	96_α	96_p	94_p
72	102_α	$93_\alpha, 98_\alpha$	$91_\alpha, 96_\alpha$	$95_\alpha, 100_\alpha$	96_α	$89_\alpha, 94_\alpha$	100_α	96_α	94_α
73	104_α	$95_p,$ $100_\alpha,$ 102_α	95_p	$97_\alpha, 99_\alpha$		$98_p,$ $100_{p,\alpha},$ $102_{p,\alpha}$	102_α	$98_{p,\alpha}$	98_p
74	99_α	102_α	$93_{p,\alpha}, 98_\alpha$	$101_\alpha,$ 104_α	$95_\alpha, 97_\alpha,$ 100_α	$95_{p,\alpha},$ 100_α	104_α	102_α	98_α
75	$101_\alpha,$ 106_α	$99_{p,\alpha}$	99_p		$97_p, 104_\alpha$	$97_p,$ $102_{p,\alpha}$	106_α	102_α	$102_{p,\alpha}$
76	103_α	$101_\alpha,$ 106_α	$97_{p,\alpha},$ 102_α	108_α	$99_\alpha, 104_\alpha$	$99_\alpha, 104_\alpha$	106_α	104_α	102_α
77	105_α		$106_{p,\alpha}$	$103_\alpha,$ 110_α	$106_\alpha,$ 108_α	101_p	110_α	106_α	104_p
78	$109_\alpha,$ 112_α	$105_\alpha,$ 108_α		$107_\alpha,$ 110_α	103_α	$101_\alpha,$ $103_\alpha,$ 106_α	109_α	106_α	106_α
79		$107_{p,\alpha},$ 109_α	107_p	$111_\alpha,$ 112_α		$105_{p,\alpha}$	112_α	110_α	110_α
80		$107_\alpha,$ 109_α	$105_\alpha,$ 108_α			$109_\alpha,$ 110_α	110_α	110_α	108_α
81				112_α	$109_{p,\alpha},$ 112_p	$109_p,$ $110_p,$ 112_p	112_α	110_p	110_p
82	111_α	$109_\alpha,$ $112_\alpha,$ 113_α	$107_{p,\alpha},$ 110_α	118_α	$105_{p,\alpha},$ 107_α	$105_p,$ $109_\alpha,$ 113_α	114_α	113_α	110_α